

Question 1

- a) We saw that the number density of non-relativistic particles with zero chemical potential in thermal equilibrium is

$$n_{\text{NR}} = g \left(\frac{mT}{2\pi} \right)^{3/2} e^{-m/T} \quad (1)$$

For $T \ll m$, this tells us that the abundance of such particle is exponentially suppressed. What is the actual physical process that is responsible for exponentially suppressing the abundance?

- b) Assuming that the process you found in part (a) is efficient, abundances will indeed get exponentially suppressed. However, the presence of a chemical potential can change that. Show that the difference in abundances between a particle (number density n) and its antiparticle (number density $\bar{n}$) is

$$n - \bar{n} = 2g \left(\frac{mT}{2\pi} \right)^{3/2} e^{-m/T} \sinh \frac{\mu}{T} \quad (2)$$

Where μ can be a function of temperature.

- c) Near the epoch of recombination ($T \sim 1$ eV) just before neutral atoms formed, the Universe was full of non-relativistic electrons ($m_e \sim 0.5$ MeV), apparently contradicting Eq. (1) above. How is this possible? Which eventual value of the chemical potential do you need to make this possible?

- a) In the case of $\mu = 0$, the number densities for the particle and its antiparticle are equal to each other.

As such, and for practical purposes, the particle-antiparticle annihilation process can take place freely (that is, the universe doesn't have to worry about keeping a certain asymmetry between their number densities as a result of a non-zero chemical potential).

The difference between this scenario and that of higher energies (higher T) is that a low T , producing the pair of particles becomes too "expensive" for the universe to carry out, so there is a single process taking place and this process annihilates them. This gives rise to this exponential suppression.

In other words, while pair-production "shuts down," annihilation still takes place.

- b) Suppose that the chemical potential is not zero. Further, if μ_X is the chemical potential of the particle (X), then the chemical potential of the antiparticle $\bar{X}$ is

$$\mu_{\bar{X}} = -\mu_X \quad (3)$$

Recall that

$$n_{\text{NR}}(T) \equiv n(T) = g \left(\frac{mT}{2\pi} \right)^{3/2} \exp \left(-\frac{m - \mu}{T} \right) \quad (4)$$

So the number density for the antiparticle is

$$\bar{n}(T) = g \left(\frac{mT}{2\pi} \right)^{3/2} \exp \left(-\frac{m + \mu}{T} \right) \quad (5)$$

Together and with some algebraic manipulations, we observe that

$$\begin{aligned} n - \bar{n} &= g \left(\frac{mT}{2\pi} \right)^{3/2} \left(\exp \left(-\frac{m - \mu}{T} \right) - \exp \left(-\frac{m + \mu}{T} \right) \right) \\ &= g \left(\frac{mT}{2\pi} \right)^{3/2} \exp \left(-\frac{m}{T} \right) [\exp(+\mu/T) - \exp(-\mu/T)] \\ &= g \left(\frac{mT}{2\pi} \right)^{3/2} \exp \left(-\frac{m}{T} \right) 2 \frac{[\exp(+\mu/T) - \exp(-\mu/T)]}{2} \\ &= g \left(\frac{mT}{2\pi} \right)^{3/2} \exp \left(-\frac{m}{T} \right) 2 \sinh(\mu/T) \\ &= 2g \left(\frac{mT}{2\pi} \right)^{3/2} \exp \left(-\frac{m}{T} \right) \sinh(\mu/T) \end{aligned} \quad (6)$$

Where the fact that $\sin(x) = \frac{e^x - e^{-x}}{2}$ was used.

- c) Equation 1 assumes zero chemical potential. This is not the case in our universe! Suppose that $n \gg \bar{n}$ so that $n - \bar{n} \simeq n$. Then

$$n \simeq n - \bar{n} = 2g \left(\frac{mT}{2\pi} \right)^{3/2} \exp \left(-\frac{m}{T} \right) \sinh(\mu/T) \quad (7)$$

Observe that $m \gg T$. In this case, we must have

$$\mu \simeq m = 0.5 \text{ MeV} \quad (8)$$

To make the equations balance out correctly.