## ST: General Relativity Extra Problem 8

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Q1. Consider the Poincaré half-plane, which has for metric

$$ds^2 = \frac{a^2}{y^2} \left( dx^2 + dy^2 \right),$$

with y > 0, and where a is a constant.

(a). Compute the length of a x = constant line segment between the coordinates  $y_1$  and  $y_2$ , with  $y_2 > y_1$ . Could an observer reach y = 0 by traveling a finite distance.

(b). Show that the geodesics in this space are either semi-circles with centers located on the x-axis or x =constant lines.

(c). Is this space curved? Is this a maximally symmetric space?.

Sol.

We are given that

$$g_{\mu\nu} = \begin{pmatrix} \frac{a^2}{y^2} & \\ & \frac{a^2}{y^2} \\ & \frac{a^2}{y^2} \end{pmatrix} \quad \& \quad g^{\mu\nu} = \begin{pmatrix} \frac{y^2}{a^2} & \\ & \frac{y^2}{a^2} \\ & \frac{y^2}{a^2} \end{pmatrix}$$

We list all non-zero derivatives of metric tensor

$$\partial_y g_{xx} = -\frac{2a^2}{y^3}, \quad \partial_y g_{yy} = -\frac{2a^2}{y^3}.$$

Let's compute Christoffel symbols first

$$\Gamma^{\alpha}_{\mu\nu} = \frac{1}{2} g^{\alpha\beta} \left( \partial_{\mu} g_{\nu\beta} + \partial_{\nu} g_{\beta\mu} - \partial_{\beta} g_{\mu\nu} \right),$$

setting  $\alpha = x$  one has

$$\Gamma^x_{\mu\nu} = \frac{1}{2}g^{xx}\left(\partial_\mu g_{\nu x} + \partial_\nu g_{x\mu} - \partial_x g_{\mu\nu}\right)$$

the only non-zero terms are  $\mu=y,\;\nu=x$  or  $\mu=x,\;\nu=y$ 

$$\Gamma_{xy}^{x} = \Gamma_{yx}^{x} = \frac{1}{2}g^{xx}\left(\partial_{y}g_{xx}\right) = -\frac{1}{y},$$

likewise, setting  $\alpha = y$  we have

$$\Gamma^{y}_{\mu\nu} = \frac{1}{2} g^{yy} \left( \partial_{\mu}g_{\nu y} + \partial_{\nu}g_{y\mu} - \partial_{y}g_{\mu\nu} \right)$$

the only non-zero terms are  $\mu = x$ ,  $\nu = x$  and  $\mu = y$ ,  $\nu = y$ 

$$\Gamma_{xx}^{y} = \frac{1}{2}g^{yy} \left(-\partial_{y}g_{xx}\right) = \frac{1}{y} \quad \& \quad \Gamma_{yy}^{y} = \frac{1}{2}g^{yy} \left(\partial_{y}g_{yy}\right) = -\frac{1}{y}.$$

Geodesic equation is given by

$$\frac{d^2x^{\mu}}{ds^2} + \Gamma^{\mu}_{\alpha\beta} \frac{dx^{\alpha}}{ds} \frac{dx^{\beta}}{ds} = 0,$$

**Caveat**! Since we are not in Minkowski space, thus we don't have such relation  $d\tau = -ds^2$ , thus instead using proper time  $d\tau$  as variable for geodesic equation, it would be more appropriate to use ds as variable. Hence

$$\int \frac{d^2x}{ds^2} - \frac{2}{y}\frac{dx}{ds}\frac{dy}{ds} = 0 \tag{1}$$

$$\left(\frac{d^2y}{ds^2} + \frac{1}{y}\left(\left(\frac{dx}{ds}\right)^2 - \left(\frac{dy}{ds}\right)^2\right) = 0$$
(2)

let's focus on eq (1), one may realize that integral factor is  $y^{-2}$ , thus eq (1) read

$$\frac{d}{ds}\left(\frac{dx}{ds}y^{-2}\right) = 0 \implies \frac{dx}{ds} = Cy^2, \ C \in \mathbb{R}.$$

(a). Now if C = 0, we get x = constant, in this case geodesic is simplify vertical line, thus

$$\ell = \int_{y_1}^{y_2} \sqrt{ds^2} = \int_{y_1}^{y_2} \frac{a}{y} \sqrt{\left(\frac{dx}{ds}\frac{ds}{dy}\right)^2 + 1} \, dy = a \int_{y_1}^{y_2} \frac{dy}{y} = a \log \frac{y_2}{y_1}$$

i.e., the length of x = constant line segment between  $y_1$  and  $y_2$  is given by  $a \log \frac{y_2}{y_1}$ , notice that observer would never reach y = 0 by traveling finite distance.

(b). We have shown in part (a) that if C = 0 geodesics is just vertical lines, next assume  $C \neq 0$ , then consider following equations

$$\left(\frac{dx}{ds} = Cy^2\right) \tag{3}$$

$$\left(\frac{a^2}{y^2}\left(\left(\frac{dx}{ds}\right)^2 + \left(\frac{dy}{ds}\right)^2\right) = 1$$
(4)

use the fact that  $\frac{dy}{ds} = \frac{dy}{dx}\frac{dx}{ds}$ , we may solve  $\frac{dy}{dx}$  from equation (3) and (4),

$$\begin{aligned} \frac{dy}{dx} &= \sqrt{\frac{y^2 - a^2 C^2 y^4}{a^2 C^2 y^4}} = \sqrt{\frac{1 - a^2 C^2 y^2}{a^2 C^2 y^2}}\\ \frac{a C y \, dy}{\sqrt{1 - a^2 C^2 y^2}} &= dx \end{aligned}$$

integrate on both sides yields

$$-\frac{\sqrt{1-a^2C^2y^2}}{aC} + a = x, \quad a \in \mathbb{R}$$
$$\frac{1}{a^2C^2} \left(1-a^2C^2y^2\right) = (x-a)^2$$
$$(x-a)^2 + y^2 = \frac{1}{a^2C^2}, \quad a, C \in \mathbb{R} \setminus \{0\}$$

as one may realize, this is semi-circles with centers located on the x-axis, as we see in following picture.



Figure 1: Geodesic in Poincaré Half-plane.

(c). Let's consider Riemann tensor, notice that the only non-zero term would be xyxy (or yxyx),

$$R^{\alpha}{}_{\beta\mu\nu} = \partial_{\mu}\Gamma^{\alpha}_{\beta\nu} - \partial_{\nu}\Gamma^{\alpha}_{\beta\mu} + \Gamma^{\alpha}_{\mu\zeta}\Gamma^{\zeta}_{\beta\nu} - \Gamma^{\alpha}_{\nu\zeta}\Gamma^{\zeta}_{\beta\mu}$$

thus

$$\begin{aligned} R^{x}_{yxy} &= \partial_{x}\Gamma^{x}_{yy} - \partial_{y}\Gamma^{x}_{yx} + \Gamma^{x}_{x\zeta}\Gamma^{\zeta}_{yy} - \Gamma^{x}_{y\zeta}\Gamma^{\zeta}_{yx} \\ &= -\partial_{y}\Gamma^{x}_{yx} + \Gamma^{x}_{xy}\Gamma^{y}_{yy} - \Gamma^{x}_{yx}\Gamma^{x}_{yx} \\ &= -\frac{1}{y^{2}} + \frac{1}{y^{2}} - \frac{1}{y^{2}} \\ &= -\frac{1}{y^{2}} \end{aligned}$$
$$\begin{aligned} R^{y}_{xyx} &= g^{yy}g_{xx}R^{x}_{yxy} \\ &= -\frac{1}{y^{2}} \end{aligned}$$

notice that  $R^x_{yyx} = -R^x_{yxy}$  and  $R^y_{xxy} = -R^y_{xyx}$ , in fully covariant form, non-zero components are

$$R_{xyxy} = R_{yxyx} = -\frac{a^2}{y^4} \quad \& \quad R_{xyyx} = R_{yxxy} = \frac{a^2}{y^4}$$

Ricci tensor is given by

$$R_{\mu\nu} = R^{\alpha}{}_{\mu\alpha\nu}$$

all that survives are two components

$$R_{xx} = R_{yy} = -\frac{1}{y^2},$$

finally Ricci scalar

$$R = g^{\mu\nu} R_{\mu\nu} = -\frac{2}{a^2}.$$

The space has negative curvature, hence it is a *hyperbolic* like space. Finally, Maximally symmetric space is a space that is both homogeneous and isotropic, such space has largest number of Killing vector fields given by  $\frac{N(N+1)}{2}$ , where  $N = \dim M$  is dimension of our manifold, further the following conditions would hold for such space

Ricci scalar R is a constant.

Ricci tensor is proportional to the metric tensor,  $R_{\mu\nu} = \frac{R}{N}g_{\mu\nu}$ . Riemann curvature tensor is given by  $R_{\alpha\beta\mu\nu} = \frac{R}{N(N-1)} \left(g_{\alpha\mu}g_{\beta\nu} - g_{\alpha\nu}g_{\beta\mu}\right)$ .

It is easy to see that Poincaré half-plane is maximally symmetric space.