

# PHYS 480/581 General Relativity

## Extra Problems #11

### Question 1.

---

Let us derive the Tolman-Oppenheimer-Volkoff equation, which is the General relativistic version of the hydrostatic equilibrium equation. To do so, we need to find a solution to the Einstein that is valid *inside* a static, spherically symmetry, non-rotating objects like a star. This is thus not a vacuum solution, and we need to solve for the full Einstein equation

$$R_{\mu\nu} = 8\pi G(T_{\mu\nu} - \frac{1}{2}g_{\mu\nu}T). \quad (1)$$

Using the same arguments as for the Schwarzschild solution, our trial metric will admit the diagonal form

$$ds^2 = -A(r,t)dt^2 + B(r,t)dr^2 + r^2(d\theta^2 + \sin^2\theta d\phi^2). \quad (2)$$

We have already shown in class that the source term for a perfect fluid admits the form

$$T_{\mu\nu} - \frac{1}{2}g_{\mu\nu}T = (\rho + p)u_\mu u_\nu + \frac{1}{2}(\rho - p)g_{\mu\nu}, \quad (3)$$

where  $\rho$  is the energy density,  $p$  is the pressure, and  $u^\mu$  is the 4-velocity of the fluid. The fluid forming the object is at rest, and thus the spatial components of the fluid 4-velocity vanish. Spherical symmetry implies that  $\rho = \rho(r)$  and  $p = p(r)$ . From the normalization of the 4-velocity, we have  $u_t = -\sqrt{A}$ .

(a) Show that the Einstein equation implies

$$R_{tt} = 4\pi G(\rho + 3p)A, \quad (4)$$

$$R_{rr} = 4\pi G(\rho - p)B, \quad (5)$$

$$R_{\theta\theta} = 4\pi G(\rho - p)r^2, \quad (6)$$

and thus that

$$\frac{R_{tt}}{A} + \frac{R_{rr}}{B} + 2\frac{R_{\theta\theta}}{r^2} = 16\pi G\rho. \quad (7)$$

(b) Use the Ricci tensor components for the above trial metric to show that the above equation reduces to

$$\frac{d}{dr} \left[ r \left( 1 - \frac{1}{B} \right) \right] = 8\pi G\rho r^2. \quad (8)$$

(c) Integrate both sides to get

$$B(r) = \left[ 1 - \frac{2Gm(r)}{r} \right]^{-1}, \quad (9)$$

where

$$m(r) \equiv \int_0^r 4\pi\rho r^2 dr. \quad (10)$$

Is  $m(r)$  the mass enclosed within a radius  $r$ ?

- (d) To find  $A(r)$ , it is easiest to use energy conservation  $\nabla_\nu T^{\mu\nu} = 0$ . Show that the  $\mu = r$  component of this equation implies

$$\frac{1}{A} \frac{dA}{dr} = -\frac{2}{\rho + p} \frac{dp}{dr}. \quad (11)$$

- (e) Derive the Tolman-Oppenheimer-Volkoff equation by solving the above for  $dp/dr$  and using the  $\theta\theta$  component of the Einstein equation to solve for  $A$ .

$$\frac{dp}{dr} = -\frac{\rho + p}{r^2} \left[ \frac{4\pi Gpr^3 + Gm(r)}{1 - 2Gm(r)/r} \right]. \quad (12)$$